

R-matrix analysis of reactions in the ${}^9\text{B}$ compound system

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National Lab**

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Outline

- **Motivation: ENDF/ENSDF evaluations, ${}^9\text{B}$ in BBN/ ${}^7\text{Li}$ destruction**
- **R-matrix formalism: T-matrix/observables, EM channels, EDA code**
- **Summary of ${}^9\text{B}$ data: DCS, σ**
- **Analysis: χ^2/N_{data} , resolution broadening**
 - Used earlier 3-channel evaluation by G. Hale, added capture channel
- **Resonance structure: implications for BBN**
- **Summary, findings & future work**

Motivation

■ Cross section evaluation & resonance structure

→ *Nucl. Phys. A745, 155, 2004(2011)*

E_x^a (MeV \pm keV)	$J^\pi; T$	$\Gamma_{c.m.}$ (keV)	Decay
16.024 \pm 25	$T = (\frac{1}{2})$	180 \pm 16	
16.71 \pm 100 ^h	$(\frac{5}{2}^+); (\frac{1}{2})$		
17.076 \pm 4	$\frac{1}{2}^-; \frac{3}{2}$	22 \pm 5	$(\gamma, {}^3\text{He})$
17.190 \pm 25		120 \pm 40	p, d, ${}^3\text{He}$
17.54 \pm 100 ^{h,i}	$(\frac{7}{2}^+); (\frac{1}{2})$		
17.637 \pm 10 ⁱ		71 \pm 8	p, d, ${}^3\text{He}, \alpha$

■ Astrophysical applications

→ Big bang nucleosynthesis

- Nuclear physics solution to ${}^7\text{Li}$ predicted overproduction problem? (*cf. Hoyle*)
- Details next slide.

■ Purpose within Los Alamos Nat. Lab programmatic

→ Continue the R-matrix program for various end-users

→ Ongoing/upcoming analysis releases: ${}^7\text{Be}$, ${}^{13}\text{C}$ [*G. Hale Tues. Session GA 2*], ${}^{14}\text{C}$, ${}^{17}\text{O}$, ...

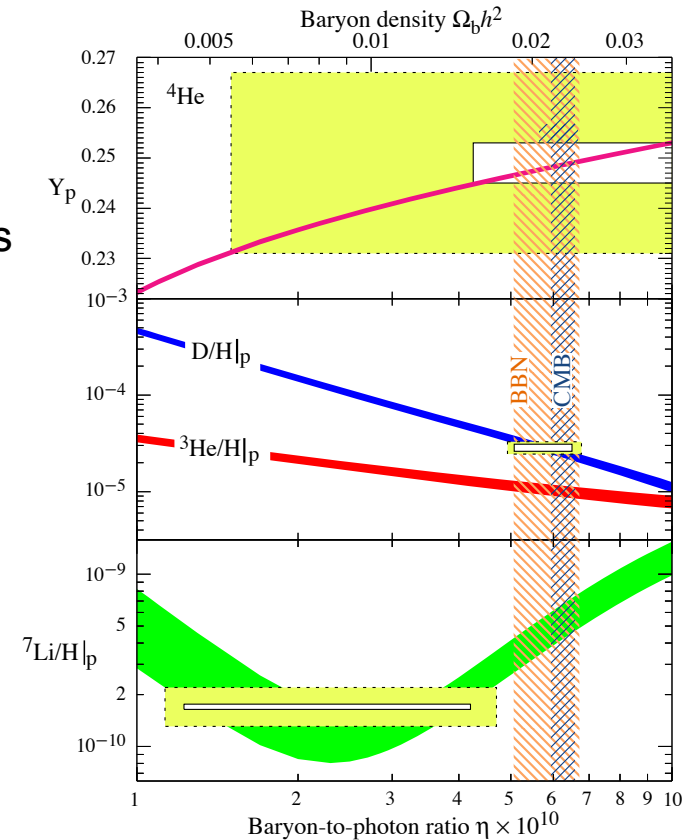
A nuclear physics solution to the BBN ${}^7\text{Li}$ problem?

■ Primordial nucleosynthesis

- Probes early universe w/in standard model
- Big-bang nucleosynthesis: D , ${}^4\text{He}$, ${}^7\text{Li}$ abundances
- D , ${}^4\text{He}$ abundances agree with theo/expl uncertainties
- At η_{Wmap} (CMB) ${}^7\text{Li}/\text{H}|_{\text{BBN}} \sim (2.2-4.2) \cdot {}^7\text{Li}/\text{H}|_{\text{halo}}$ *
- Discrepancy $\sim 4.5-5.5\sigma$ → the “Li problem”

■ Resonant destruction ${}^7\text{Li}$ (Hoyle-type solution)

- Prod. mass 7 “well understood”; destruction not
- Cyburt & Pospelov *arXiv:0906.4373; IJMPE, 21(2012)*
 - ${}^7\text{Be}(d,p)\alpha$ & ${}^7\text{Be}(d,\gamma){}^9\text{B}$ resonant enhancement
 - Identify ${}^9\text{B}$ $E_{5/2^+} \approx 16.7 \text{ MeV} \approx E_{\text{thr}}(d+{}^7\text{Be}) + 200 \text{ keV}$
 - *Near threshold*
 - $(E_r, \Gamma_d) \approx (170-220, 10-40) \text{ keV}$ solve Li problem
- Chakraborty, Fields & Olive *PRD83, 063006 (2011)*
 - More general approach: $A=8,9,10$ & 11
 - Identify as possibly important: ${}^9\text{B}$, ${}^{10}\text{B}$, ${}^{10}\text{C}$
- ‘Large’ widths
 - Both conclude “large channel radius” required



NB: both approaches assume validity of TUNL-NDG tables

R-matrix formalism

INTERIOR (Many-Body) REGION
(Microscopic Calculations)

ASYMPTOTIC REGION
(S-matrix, phase shifts, etc.)

$H + \mathcal{L}_B$
compact, hermitian operator with real, discrete spectrum; eigenfunctions in Hilbert space



$$|\psi^+\rangle = (H + \mathcal{L}_B - E)^{-1} \mathcal{L}_B |\psi^+\rangle$$

SURFACE

$$\mathcal{L}_B = \sum_c |c\rangle \langle c| \left(\frac{\partial}{\partial r_c} r_c - B_c \right)$$

$$\langle \mathbf{r}_c | c \rangle = \frac{\hbar}{\sqrt{2\mu_c a_c}} \frac{\delta(r_c - a_c)}{r_c} [(\phi_n^{\mu_1} \otimes \phi_n^{\mu_2})_s^{\mu} \otimes Y_l^m(\hat{\mathbf{r}}_c)]_j^M$$

$$R_{c'c} = \langle c' | (H + \mathcal{L}_B - E)^{-1} | c \rangle = \sum_{\lambda} \frac{\langle c' | \lambda \rangle \langle \lambda | c \rangle}{E_{\lambda} - E}$$

$$\langle r_{c'} | \psi_c^+ \rangle = -I_{c'}(r_{c'}) \delta_{c'c} + O_{c'}(r_{c'}) S_{c'c}$$

or equivalently,

$$\langle r_{c'} | \psi_c^+ \rangle = F_{c'}(r_{c'}) \delta_{c'c} + O_{c'}(r_{c'}) T_{c'c}$$

$\phi_{s_i}^{\mu_i}$

nucleus internal w.f.'s

- Bloch operator $\mathcal{L}_B = \sum_c |c\rangle \langle c| \left[\frac{\partial}{\partial r_c} r_c - B_c \right]$ ensures Hermiticity of Hamiltonian restricted to internal region

- R-matrix theory: unitary, multichannel parametrization of (not just resonance) data

- Interior/Exterior regions
 - Interior: strong interactions
 - Exterior: Coulomb/non-polarizing interactions
 - Channel surface

$$\mathcal{S}_c : r_c = a_c \quad \mathcal{S} = \sum_c \mathcal{S}_c$$

- R-matrix elements
 - Projections on channel surface functions $\langle \mathbf{r}_c | c \rangle$ of Green's function

$$G_B = [H + \mathcal{L}_B - E]^{-1}$$

- Boundary conditions

$$B_c = \frac{1}{u_c(a_c)} \frac{du_c}{dr_c} \Big|_{r_c=a_c}$$

- E-M channels
 - Next slide

Electromagnetic channels

■ One-photon sector of Fock space

→ Photon 'wave function'

$$\mathbf{A}_{\mathbf{k}}(\mathbf{r}) = \left(\frac{2}{\pi \hbar c} \right)^{1/2} \sum_{jm} i^j \sum_{\lambda', \lambda=e, m, 0} \mathbf{Y}_{jm}^{(\lambda')}(\hat{\mathbf{r}}) u_{\lambda', \lambda}^j(r) \mathbf{Y}_{jm}^{(\lambda)}(\hat{\mathbf{k}}) \cdot \chi$$

→ Radial part

$$\begin{aligned} u_{ee}^j &= -[f_j'(\rho) + t_{ee}^j h_j^+(\rho)] & u_{0e}^j &= -\frac{\sqrt{j(j+1)}}{\rho} [f_j(\rho) + t_{e0}^j h_j^+(\rho)] \\ u_{mm}^j &= [f_j(\rho) + t_{mm}^j h_j^+(\rho)] & u_{0m}^j &= u_{me}^j = u_{em}^j = 0 \end{aligned}$$

→ Photon channel surface functions

$$(\mathbf{r}_c | c) = \left(\frac{\hbar c}{2\rho_\gamma} \right)^{1/2} \frac{\delta(r_\gamma - a_\gamma)}{r_\gamma} \left[\phi_{s\nu} \otimes \mathbf{Y}_{jm}^{(e,m)}(\hat{\mathbf{r}}_\gamma) \right]_{JM}$$

• Photon 'mass': $\hbar k_\gamma / c$

→ R-matrix definition preserved

$$(c' | \psi) = \sum_c R_{c'c}^B (c | \frac{\partial}{\partial r_c} r_c - B_c | \psi)$$

$$\begin{aligned} \mathbf{T} &= \rho^{1/2} \mathbf{O}^{-1} \mathbf{R}_L \mathbf{O}^{-1} \rho^{1/2} - \mathbf{F} \mathbf{O}^{-1} \\ \mathbf{R}_L &= [\mathbf{R}_B^{-1} - \mathbf{L} + \mathbf{B}]^{-1} \\ \mathbf{L} &= \rho \mathbf{O}' \mathbf{O}^{-1} \\ F &= \text{Im } \mathbf{O} \end{aligned}$$

Implementation in EDA

- **EDA = Energy Dependent Analysis**

- Adjust E_λ & $\gamma_{c\lambda}$

- **Any number of two-body channels**

- Arbitrary spins, masses, charges (incl. mass zero)

- **Scattering observables**

- Wolfenstein trace formalism

- **Data**

- Normalization

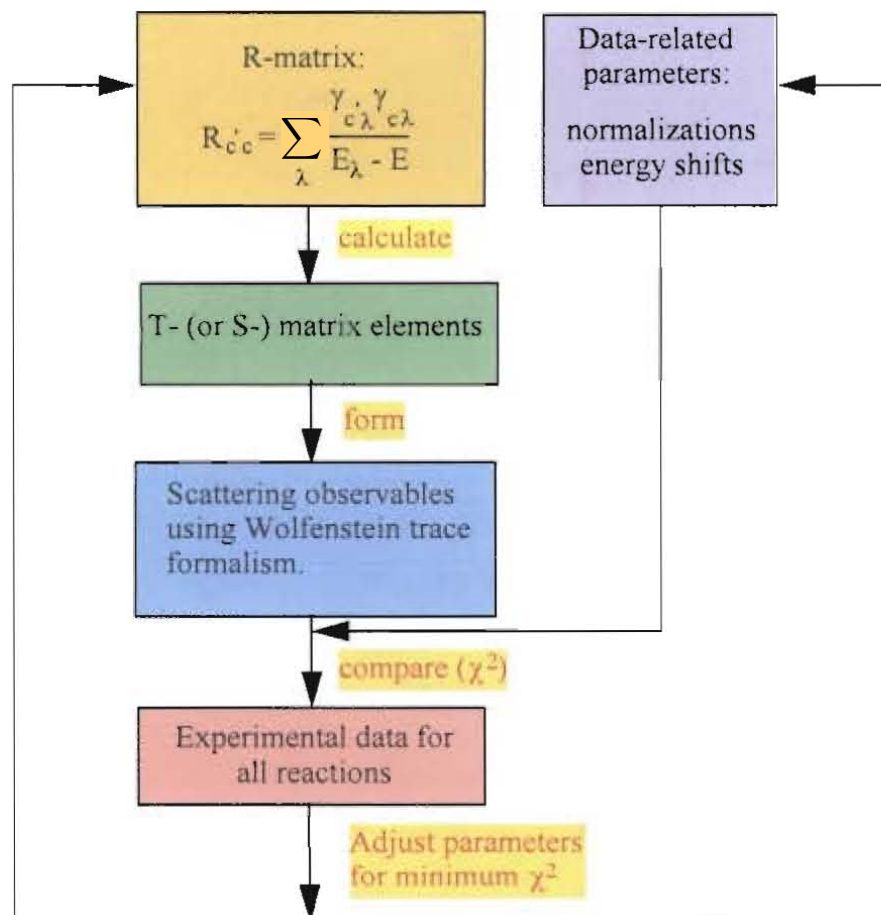
- Energy shifts

- Energy resolution/spread

- **Fit solution**

$$\chi_{EDA}^2 = \sum_i \left[\frac{nX_i(\mathbf{p}) - R_i}{\delta R_i} \right]^2 + \left[\frac{nS - 1}{\delta S/S} \right]^2$$

- **Covariance determined**



Summary of included ${}^9\text{B}$ data

- **${}^6\text{Li}+{}^3\text{He}$ elastic** *Buzhinski et.al., Izv. Rossiiskoi Akademii Nauk, Ser.Fiz., Vol.43, p.158 (1979)*
 - Differential cross section
 - $1.30 \text{ MeV} < E({}^3\text{He}) < 1.97 \text{ MeV}$
- **${}^6\text{Li}+{}^3\text{He} \rightarrow \text{p}+{}^8\text{Be}^*$** *Elwyn et.al., Phys. Rev. C 22, 1406 (1980)*
 - Integrated cross section
 - Quasi-two-body, excited-state averaged final channel
 - $0.66 \text{ MeV} < E({}^3\text{He}) < 5.00 \text{ MeV}$
- **${}^6\text{Li}+{}^3\text{He} \rightarrow \text{d}+{}^7\text{Be}$** *D.W. Barr & J.S. Gilmore, unpublished (1965)*
 - Integrated cross section
 - $0.42 \text{ MeV} < E({}^3\text{He}) < 4.94 \text{ MeV}$
- **${}^6\text{Li}+{}^3\text{He} \rightarrow \gamma+{}^9\text{B}$** *Aleksic & Popic, Fizika 10, 273-278 (1978)*
 - Integrated cross section
 - $0.7 \text{ MeV} < E({}^3\text{He}) < 0.825 \text{ MeV}$
 - New to ${}^9\text{B}$ analysis
- **Data for future evaluation**
 - Separate ${}^8\text{Be}^*$ states
 - 2^+ @200 keV [16.9 MeV], 1^+ @650 keV [17.6 MeV], 1^+ @1.1 MeV [18.2 MeV]
 - $\text{n}+{}^8\text{B}$: $E_{\text{thresh}}({}^3\text{He}) = 3 \text{ MeV}$
 - Simultaneous analysis with ${}^9\text{Be}$ mirror system

All data from
EXFOR/CSISRS
database (in
C4 format)

R-matrix configuration in EDA code

Hadronic channels (in blue, not included)

$A_1 A_2 \pi$	${}^3\text{He}{}^6\text{Li}^+(1)$		$p{}^8\text{Be}^{*+}(2)$		$d{}^7\text{Be}^-(3)$		
S	$\frac{3}{2}$	$\frac{1}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{1}{2}$
ℓ							
0	${}^4S_{3/2}$	${}^2S_{1/2}$	${}^6S_{5/2}$	${}^4S_{3/2}$	${}^6S_{5/2}$	${}^4S_{3/2}$	${}^2S_{1/2}$
1	${}^4P_{5/2,3/2,1/2}$	${}^2P_{3/2,1/2}$	${}^6P_{7/2,5/2,3/2}$	${}^4P_{5/2,3/2,1/2}$	${}^6P_{7/2,5/2,3/2}$	${}^4P_{5/2,3/2,1/2}$	${}^2P_{3/2,1/2}$
2	${}^4D_{7/2,5/2,3/2,1/2}$	${}^2D_{5/2,3/2}$	${}^6D_{9/2,7/2,5/2,3/2,1/2}$	${}^4D_{7/2,5/2,3/2,1/2}$	${}^6D_{9/2,7/2,5/2,3/2,1/2}$	${}^4D_{7/2,5/2,3/2,1/2}$	${}^2D_{5/2,3/2}$
$E_{\text{thr}}(\text{CM, MeV})$	16.6		16.7		16.5		

Electromagnetic channel: $\gamma + {}^9\text{B} \rightarrow E_1^{3/2}, M_1^{5/2}, M_1^{3/2}, M_1^{1/2}, E_1^{5/2}, E_1^{1/2}$

Full model space:
state number;
channel pair;
LS; J; channel
radius [fm]

1	1 4s 3/2	7.50000000f	20	1 4p 1/2	7.50000000f
2	1 4d 3/2	7.50000000f	21	1 2p 1/2	7.50000000f
3	1 2d 3/2	7.50000000f	22	2 4p 1/2	5.50000000f
4	2 4s 3/2	5.50000000f	23	3 2s 1/2	7.00000000f
5	3 6p 3/2	7.00000000f	24	4 M1 1/2	50.00000000f
6	3 4p 3/2	7.00000000f	25	1 4d 7/2	7.50000000f
7	3 2p 3/2	7.00000000f	26	3 6p 7/2	7.00000000f
8	4 E1 3/2	50.00000000f	27	1 4d 5/2	7.50000000f
9	1 4p 5/2	7.50000000f	28	1 2d 5/2	7.50000000f
10	2 6p 5/2	5.50000000f	29	2 6s 5/2	5.50000000f
11	2 4p 5/2	5.50000000f	30	3 6p 5/2	7.00000000f
12	3 6s 5/2	7.00000000f	31	3 4p 5/2	7.00000000f
13	4 M1 5/2	50.00000000f	32	4 E1 5/2	50.00000000f
14	1 4p 3/2	7.50000000f	33	1 4d 1/2	7.50000000f
15	1 2p 3/2	7.50000000f	34	1 2s 1/2	7.50000000f
16	2 6p 3/2	5.50000000f	35	3 4p 1/2	7.00000000f
17	2 4p 3/2	5.50000000f	36	3 2p 1/2	7.00000000f
18	3 4s 3/2	7.00000000f	37	4 E1 1/2	50.00000000f
19	4 M1 3/2	50.00000000f	38	2 6p 7/2	5.50000000f

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Analysis result: resonance structure

Ex (MeV)	Jpi	Gamma (keV)	Er (MeV)	ImEr (MeV)	E (3He)	Strength
16.46539	1/2-	768.46	-.1369	-0.3842	-0.2054	0.06 weak
17.11317	1/2-	0.14	0.5109	-0.6771E-04	0.7664	1.00 strong
17.20115	5/2-	871.63	0.5989	-0.4358	0.8984	0.40 weak
17.28086	3/2-	147.78	0.6785	-0.0739	1.0178	0.77 strong
17.66538	5/2+	33.33	1.0631	-0.0167	1.5947	0.98 strong
17.83619	7/2+	2036.21	1.2339	-1.0181	1.8509	0.15 weak
17.84773	3/2-	42.52	1.2454	-0.0213	1.8681	0.97 strong
18.04821	3/2+	767.11	1.4459	-0.3836	2.1689	0.54 weak
18.42292	1/2+	5446.32	1.8206	-2.7232	2.7309	0.03 weak
18.67716	1/2-	10278.41	2.0749	-5.1392	3.1124	0.15 weak
19.60923	3/2-	1478.22	3.0069	-0.7391	4.5104	0.52 weak

S-matrix pole & residue *Hale, Brown, Jarmie PRL 59 '87*

$$\mathcal{E}_{\lambda'\lambda} = E_\lambda \delta_{\lambda'\lambda} - \sum_c \gamma_{c\lambda'} [L_c(E) - B_c] \gamma_{c\lambda}$$

$$E_0 = E_r - i\Gamma/2 \quad \text{residue: } i\rho_0\rho_0^T$$

**NB: no strong resonance seen
~100 keV of $^3\text{He}+^6\text{Li}$ threshold**

$$\text{strength} = \frac{1}{\Gamma} \rho_0^\dagger \rho_0 = \frac{1}{\Gamma} \sum_c \Gamma_c$$

$$\rho_{0c} = \left(\frac{2k_{0c}a_c}{N} \right)^{1/2} \mathcal{O}_c^{-1}(k_{0c}a_c) \sum_\lambda (\lambda|\mu_0)$$

$$N = \sum_{\lambda'\lambda} (\lambda|\mu_0)(\lambda'|\mu_0) \left[\delta_{\lambda'\lambda} + \sum_c \gamma_{c\lambda'} \frac{\partial L_c}{\partial E} \Big|_{E=E_0} \gamma_{c\lambda} \right]$$

$$L_c = r_c \frac{\partial \mathcal{O}_c}{\partial r_c} \mathcal{O}_c^{-1} \Big|_{r_c=a_c}$$

Analysis result: resonance structure

Ex (MeV)	Jpi	Gamma (keV)	Er (MeV)	ImEr (MeV)	E(3He)	Strength
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TUNL-NDG/ENSDF
parameters

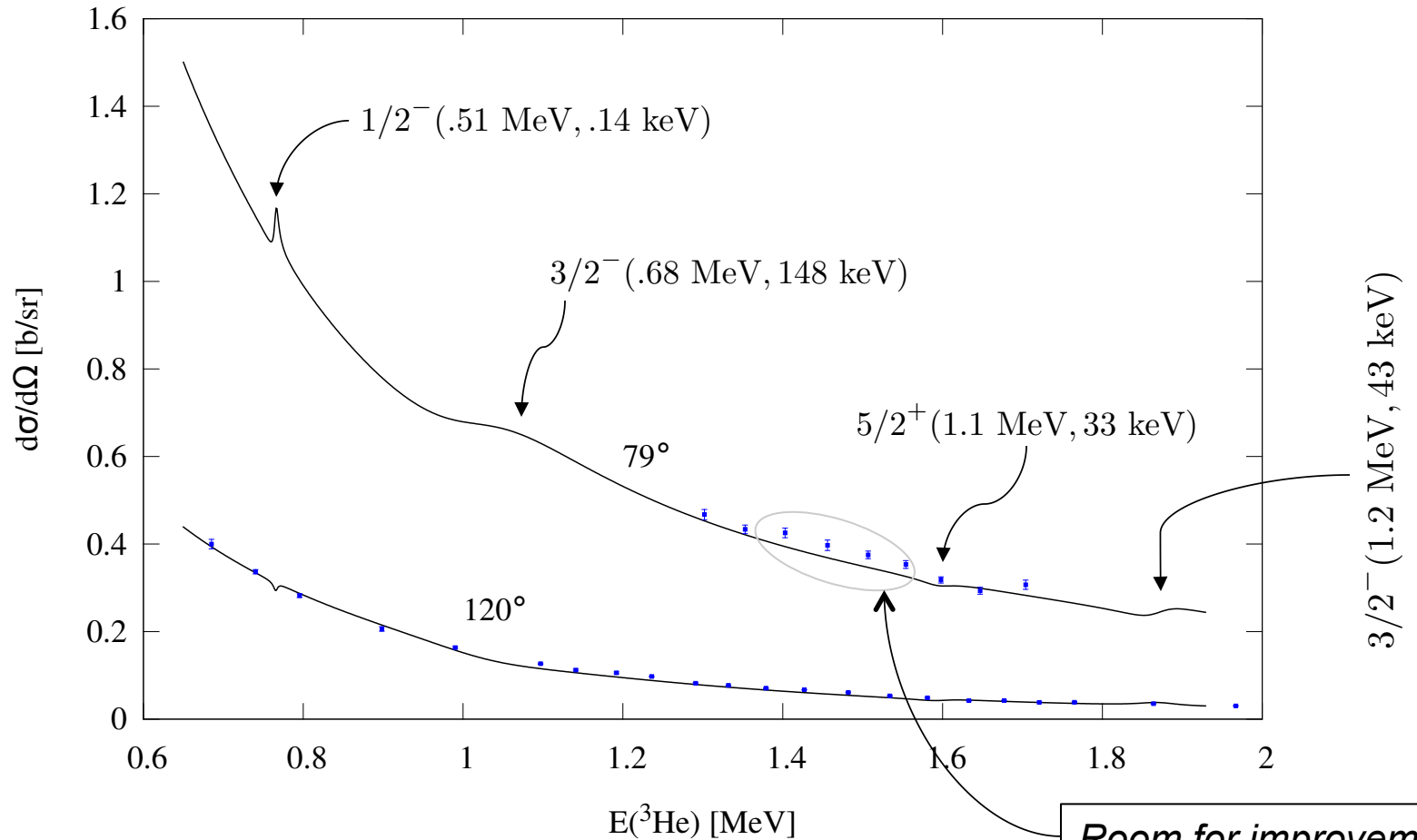
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Observable fit: ${}^3\text{He}+{}^6\text{Li}$ elastic DCS

${}^6\text{Li}({}^3\text{He},\text{Elastic})$
Differential cross section

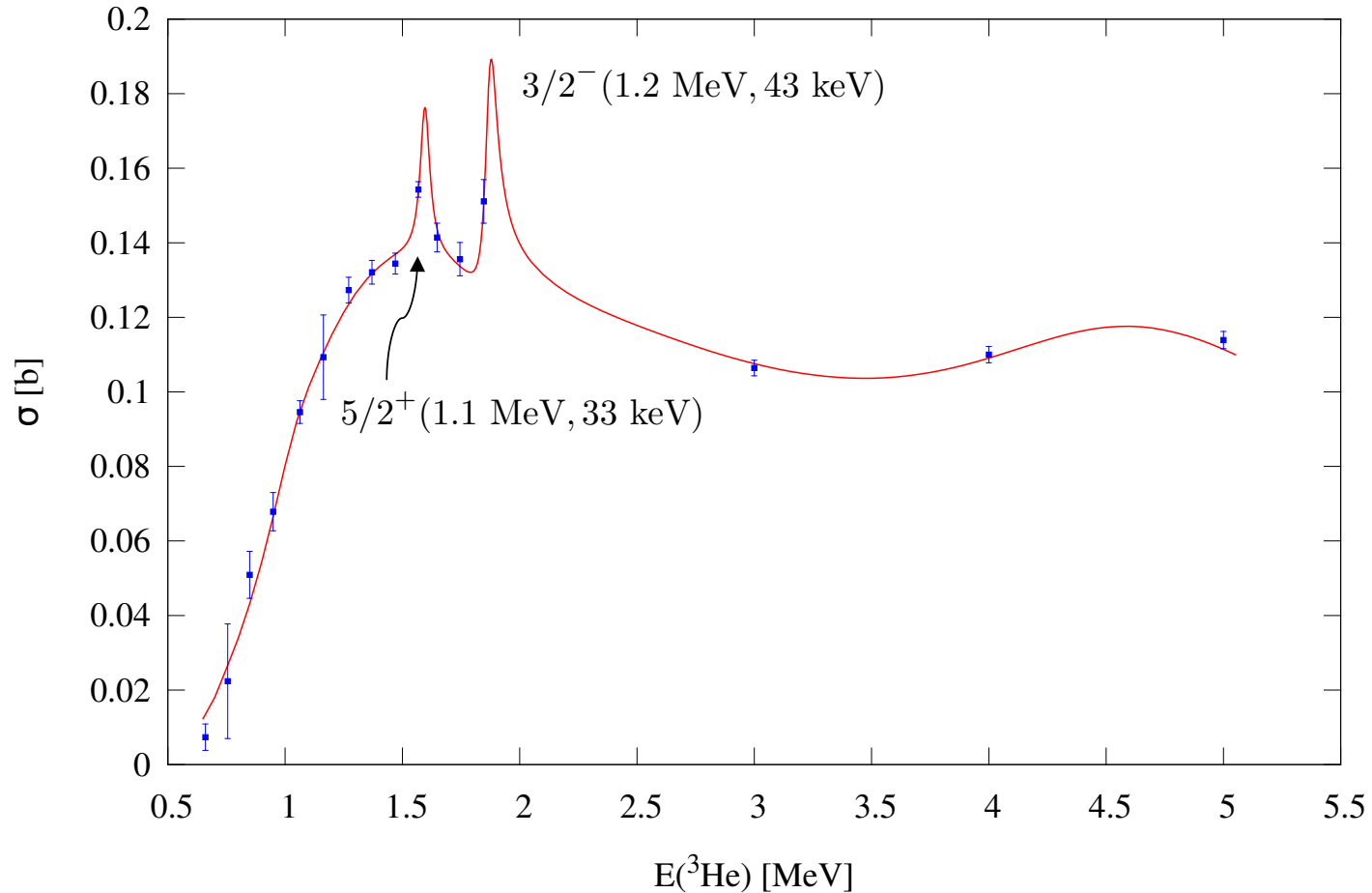
$\chi^2/N_{\text{data}} = 1.91$



Observable fit: ${}^6\text{Li}({}^3\text{He},p){}^8\text{Be}^*$ integrated x-sec

${}^6\text{Li}({}^3\text{He},p){}^8\text{Be}^*$
Integrated cross section

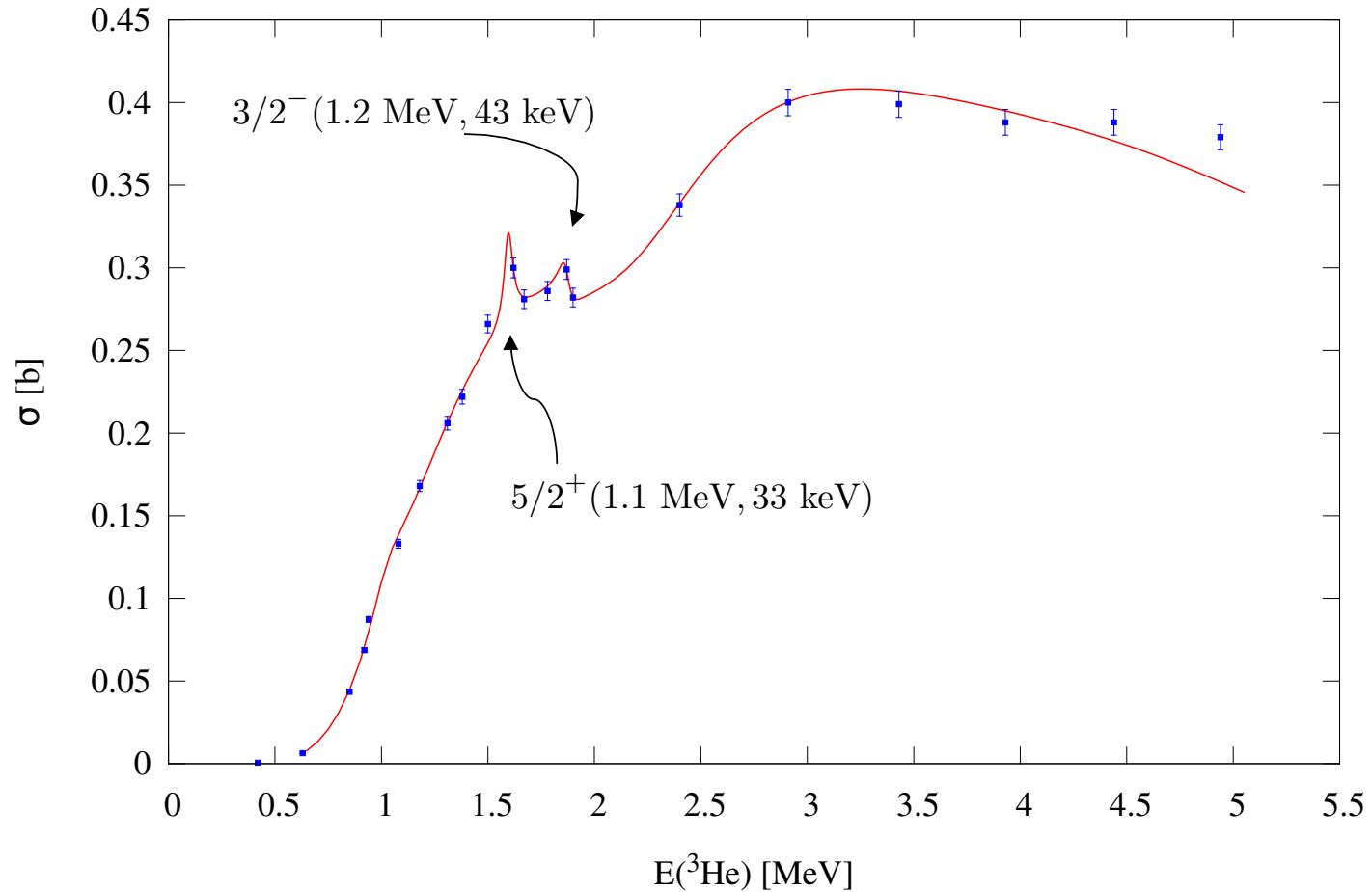
$\chi^2/N_{\text{data}} = 0.55$



Observable fit: ${}^6\text{Li}({}^3\text{He},d){}^7\text{Be}$ integrated x-sec

${}^6\text{Li}({}^3\text{He},d){}^7\text{Be}$
Integrated cross section

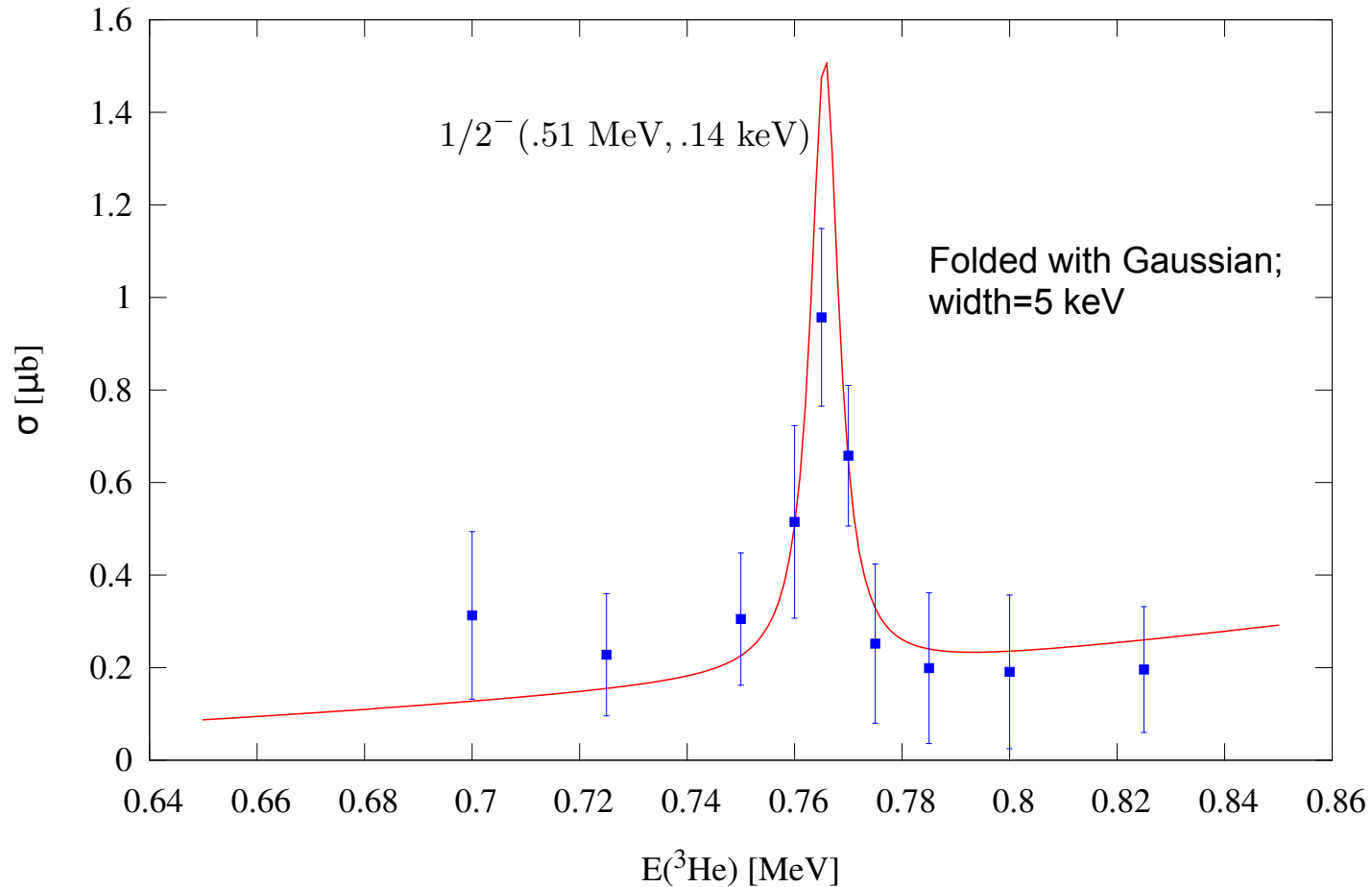
$\chi^2/N_{\text{data}} = 2.38$



Observable fit: ${}^6\text{Li}({}^3\text{He},\gamma){}^9\text{B}$ integrated x-sec

${}^6\text{Li}({}^3\text{He},\gamma){}^9\text{B}$
Integrated cross section

$\chi^2/N_{\text{data}} = 0.37$



Summary, findings & future work

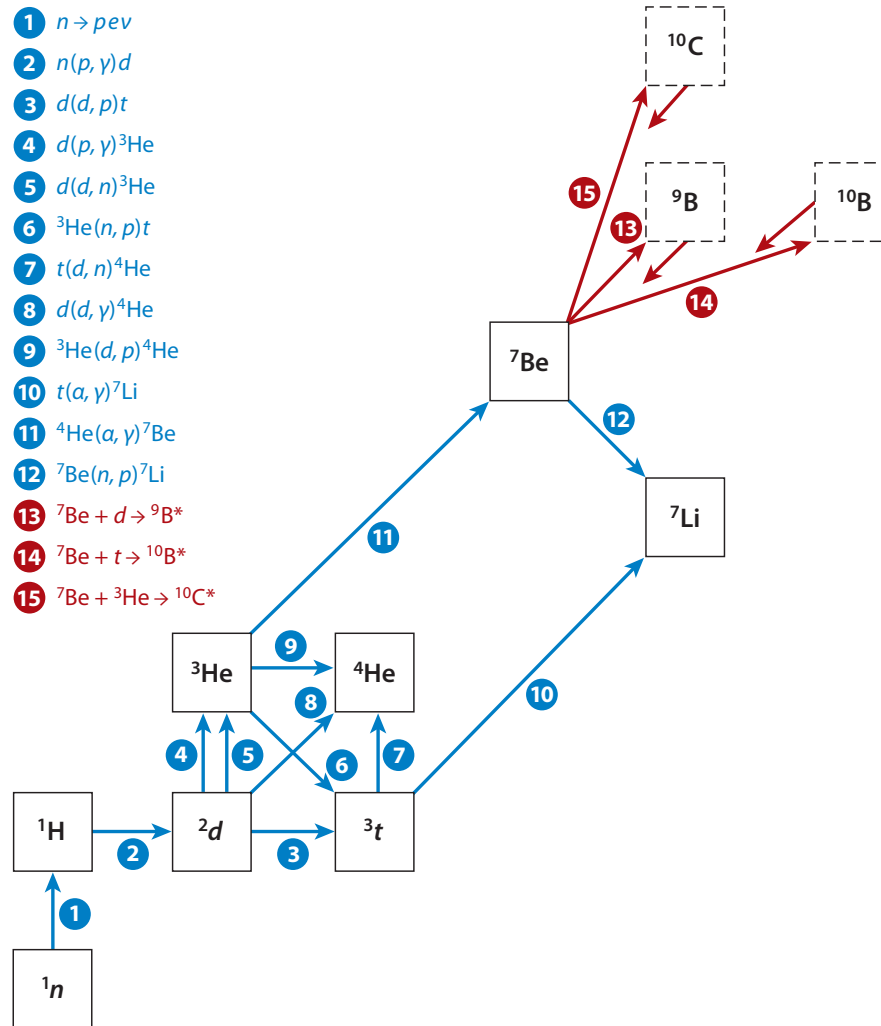
- Nuclear physics/microphysics explanations for the “Li problem” have been entertained
- There are no resonances in ${}^9\text{B}$ that reside within ~ 200 (~ 100) keV of the $d+{}^7\text{Be}$ (${}^3\text{He}+{}^6\text{Li}$) threshold with ‘large’ widths 10—40 keV
- This would appear to rule out scenarios considered by *Cybert & Pospelov (2009)* and *Chakraborty, Fields & Olive(2011)* that low-lying, robust resonance in ${}^9\text{B}$ could explain the “Li problem”
- It may be worth emphasizing that other nuclear physics explanations, such as insufficiently accurate and/or precise analyses of “known” nuclear reactions, may still be considered for the resolution of the “Li problem”
- While very useful, the TUNL-NDG/ENSDF tables may not be definitive; unitary analyses are req. and sometimes lacking; *TODO: submit new ${}^9\text{B}$ analysis*
- ***Need for dedicated, low-energy, high pol. facility***
- Improvements in the present analysis: more channels; incorporate $p+{}^8\text{Be}^*$ angular data; proper treatment three-body final states

Supplementary material

Additional slides follow

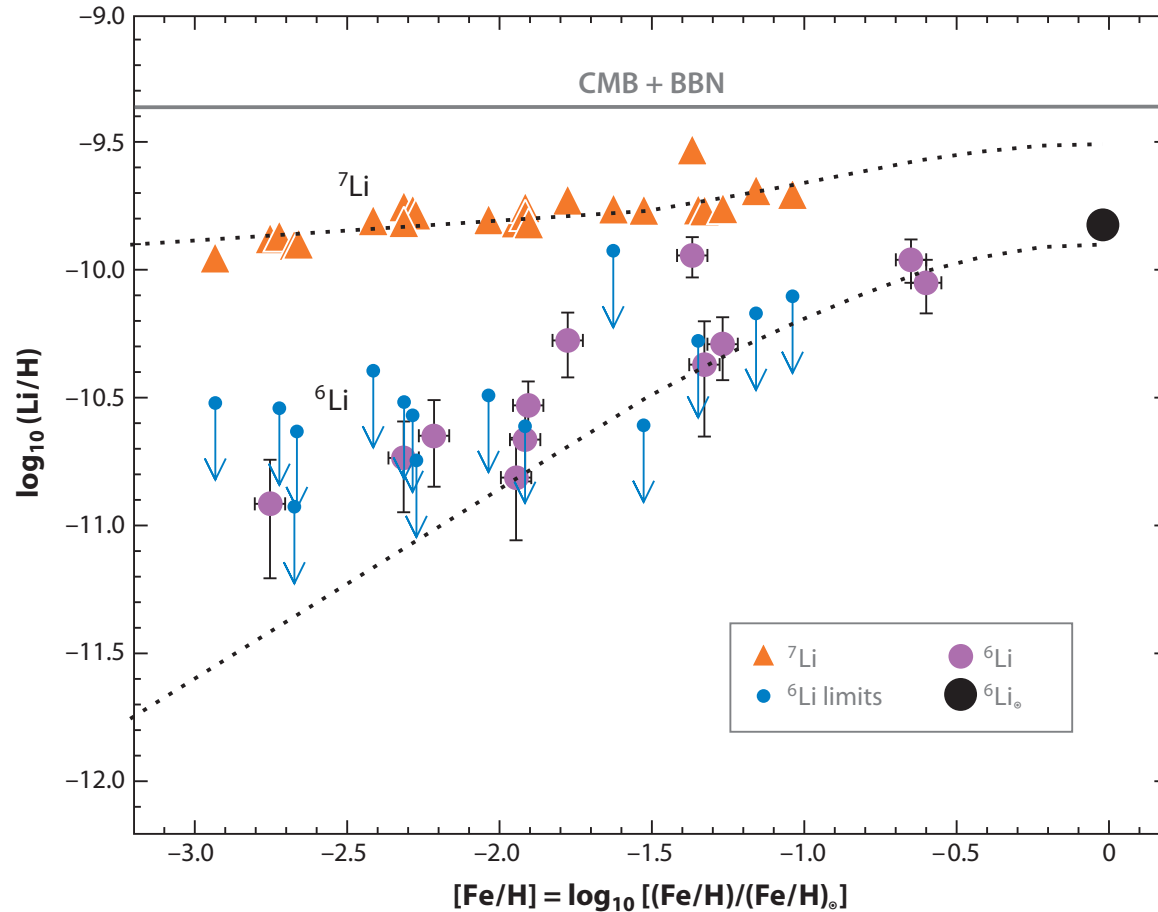
BBN reaction network (simplified)

■ **Fields** *Annu. Rev. Nucl. Part. Sci.* 2011. 61:47–68



Spite Plateau

- Measurement of primordial ${}^7\text{Li}$ from low-metallicity halo dwarf stars



Asplund M, et al. *Astrophys. J.* 644:229 (2006)